

MOMENTS OF ELLIPTIC INTEGRALS AND CRITICAL L -VALUES

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ABSTRACT. We compute the critical L -values of some weight 3, 4, or 5 modular forms, by transforming them into integrals of the complete elliptic integral K . In doing so, we prove closed form formulas for some moments of K'^3 . Many of our L -values can be expressed in terms of Gamma functions, and this also gives new lattice sum evaluations.

1. INTRODUCTION

Let K and K' denote the complete elliptic integrals of the first kind, defined as follows:

$$K(k) := \int_0^1 \frac{du}{\sqrt{(1-u^2)(1-k^2u^2)}} = \frac{\pi}{2} {}_2F_1\left(\frac{1}{2}, \frac{1}{2}; 1; k^2\right),$$
$$K'(k) := K(k'), \quad k' := \sqrt{1-k^2}.$$

One of the present authors (Wan) studied integrals of elliptic integrals, or “moments of elliptic integrals” in [14]. That paper includes several conjectures which have since been settled. The last remaining conjecture is

$$\int_0^1 K'(k)^3 dk = \frac{\Gamma^8(1/4)}{128\pi^2}. \quad (1)$$

There are several versions of (1). For instance, we can reformulate the integral using transformations for K' . We can also use integration by parts, because derivatives of elliptic integrals equal linear combinations of elliptic integrals. Some examples found in [14] include

$$\int_0^1 K'(k)^3 dk = \frac{10}{3} \int_0^1 K(k)^3 dk = 5 \int_0^1 k K'(k)^3 dk. \quad (2)$$

In the first part of this paper, we prove formula (1), and provide some intuition about how to discover related integrals. We also settle some additional integrals involving K'^3 , and give closed form evaluations of integrals containing higher powers of K' (see equations (18) and (19)). In the second part of the paper, we study a more general phenomenon, where critical L -values of *odd weight* modular forms can be

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expressed in terms of Gamma values. Some new lattice sums are produced from our investigation.

We note that Y. Zhou, in a 2013 preprint [15], used methods based on spherical harmonics to prove both equations (1) and (10). Zhou also applied his ideas to many other integrals.

2. CRITICAL L -VALUES

One of the main goals of this paper is to connect integrals like (1) to critical L -values of modular forms. We say that a function $f(\tau)$ is a modular form of weight k and level N , if it satisfies

$$f\left(\frac{a\tau + b}{c\tau + d}\right) = (c\tau + d)^k f(\tau),$$

whenever $\begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL_2(\mathbb{Z})$, and $c \equiv 0 \pmod{N}$. We also require that $f(\tau)$ be holomorphic in the upper half plane, and that it vanish at the cusps (that is, it is a cusp form). The most interesting modular forms are the Hecke eigenforms, whose Fourier series

$$f(\tau) = \sum_{n=1}^{\infty} a_n e^{2\pi i n \tau},$$

have multiplicative coefficients a_n . If we attach an L -series to $f(\tau)$,

$$L(f, s) = \sum_{n=1}^{\infty} \frac{a_n}{n^s}, \quad (3)$$

then $L(f, s)$ has a meromorphic continuation to the complex plane. We say that $L(f, j)$ is a *critical* L -value if $j \in \{1, 2, \dots, k-1\}$. These numbers typically hold arithmetic significance; some properties of critical L -values are summarized in [8]. For instance, the Birch and Swinnerton-Dyer conjecture predicts the value of $L(f, 1)$, whenever $f(\tau)$ is a weight 2 cusp form attached to an elliptic curve. In Theorem 1, we prove that equation (1) is equivalent to an explicit formula for a critical L -value of a weight 5 cusp form. In particular, we prove

$$30L(g, 4) = \int_0^1 K'(k)^3 dk, \quad (4)$$

where

$$g(\tau) := \eta^4(\tau)\eta^2(2\tau)\eta^4(4\tau),$$

and $\eta(\tau)$ is the usual Dedekind eta function. We then use properties of $g(\tau)$ to prove equation (1) in Theorem 2.

Formula (4) naturally suggests looking at critical L -values of different modular forms. Martin has classified all of the possible multiplicative eta quotients [9]. While it is possible to consider a much larger class of modular forms than just multiplicative eta quotients, these are typically the easiest examples to work with. Martin's list

contains precisely two examples of weight 5. The first is $g(\tau)$ above, and the second example is

$$h(\tau) := \frac{\eta^{38}(8\tau)}{\eta^{14}(4\tau)\eta^{14}(16\tau)}.$$

It is easy to prove that

$$192L(h, 4) = \int_0^1 \frac{K'(k)^3}{\sqrt{k}(1-k^2)^{3/4}} dk,$$

and based on the previous example, we guess that $L(h, 4)$ should also be related to Gamma values. We discovered the following identity after a brief numerical search:

$$\int_0^1 \frac{K'(k)^3}{\sqrt{k}(1-k^2)^{3/4}} dk = \frac{3\Gamma^8(1/4)}{32\sqrt{2}\pi^2}. \quad (5)$$

We started by calculating the integral on the left to high numerical precision (denoted I), and then we used the PSLQ algorithm to search for a linear dependencies in the set

$$\{\log |I|, \log \pi, \log \Gamma(1/3), \log \Gamma(1/4), \dots, \log 2, \log 3, \log 5, \dots\}.$$

Bailey and Borwein used PSLQ to discover many identities among moments of elliptic integrals – far more than what is currently proven [1].

The crucial property which allows us to prove (1) and (5), is that the attached modular forms are also binary theta functions. It is often possible to rewrite the L -functions of the modular forms as Hecke L -functions, but this connection is not always essential. In Section 4 we describe some additional integrals which arise from weight 3 cusp forms. We think that it is noteworthy that there appears to be fewer interesting formulas for L -values attached to even weight cusp forms. We discuss this in the conclusion.

3. PROOF OF THE CONJECTURES

We relate conjecture (1) to critical L -values in the following theorem. The proof is typical of the approach we use for subsequent integrals, so we spell out the details.

Theorem 1. *Let $g(\tau) = \eta^4(\tau)\eta^2(2\tau)\eta^4(4\tau)$. Formula (1) is equivalent to:*

$$L(g, 4) = \frac{\Gamma^8(1/4)}{3840\pi^2}. \quad (6)$$

Proof. The proof follows from Ramanujan-style manipulations. Set $k = \sqrt{\alpha}$, and notice that

$$\int_0^1 K'(k)^3 dk = \frac{\pi^3}{16} \int_0^1 {}_2F_1\left(\frac{1}{2}, \frac{1}{2}; 1 - \alpha\right)^3 \frac{d\alpha}{\sqrt{\alpha}}.$$

Now we make a change of variables. Set

$$q = \exp\left(-\pi \frac{{}_2F_1\left(\frac{1}{2}, \frac{1}{2}; 1 - \alpha\right)}{{}_2F_1\left(\frac{1}{2}, \frac{1}{2}; \alpha\right)}\right), \quad (7)$$

and notice that

$$d\alpha = \alpha(1-\alpha) {}_2F_1\left(\frac{1}{2}, \frac{1}{2}; \alpha\right)^2 \frac{dq}{q}.$$

It is standard to show that $q \in (0, 1)$ when $\alpha \in (0, 1)$, and q is monotone with respect to α . The integral becomes

$$\int_0^1 K'(k)^3 dk = -\frac{1}{16} \int_0^1 \sqrt{\alpha}(1-\alpha) {}_2F_1\left(\frac{1}{2}, \frac{1}{2}; \alpha\right)^5 \log^3 q \frac{dq}{q}.$$

Consider the Dedekind eta function with respect to q , where $q = e^{2\pi i\tau}$:

$$\eta(q) := q^{1/24} \prod_{n=1}^{\infty} (1 - q^n).$$

By [2, p. 124, Entry 12] we have

$$\sqrt{\alpha}(1-\alpha) {}_2F_1\left(\frac{1}{2}, \frac{1}{2}; \alpha\right)^5 = 4 \frac{\eta^{14}(q^2)}{\eta^4(q^4)},$$

and the integral reduces to

$$\int_0^1 K'(k)^3 dk = -\frac{1}{4} \int_0^1 \frac{\eta^{14}(q^2)}{\eta^4(q^4)} \log^3 q \frac{dq}{q} = -4 \int_0^1 \frac{\eta^{14}(q^4)}{\eta^4(q^8)} \log^3 q \frac{dq}{q}.$$

From [13, Entry $t_{8,12,48}$ and Entry $t_{8,18,60a}$], it is easy to prove that

$$\frac{\eta^{14}(q^4)}{\eta^4(q^8)} = 4\eta^4(q^2)\eta^2(q^4)\eta^4(q^8) + \eta^4(q)\eta^2(q^2)\eta^4(q^4).$$

Therefore the integral becomes

$$\begin{aligned} \int_0^1 K'(k)^3 dk &= -16 \int_0^1 \eta^4(q^2)\eta^2(q^4)\eta^4(q^8) \log^3 q \frac{dq}{q} - 4 \int_0^1 \eta^4(q)\eta^2(q^2)\eta^4(q^4) \log^3 q \frac{dq}{q} \\ &= -5 \int_0^1 \eta^4(q)\eta^2(q^2)\eta^4(q^4) \log^3 q \frac{dq}{q}. \end{aligned}$$

Switching to the more traditional $\eta(\tau)$ notation, we have

$$\int_0^1 K'(k)^3 dk = 30L(g, 4),$$

where $g(\tau) = \eta^4(\tau)\eta^2(2\tau)\eta^4(4\tau)$. It follows that (1) and (6) are equivalent. \square

In order to prove (4), we require the fact that g is a binary theta function. Glaisher [6] showed that

$$g(\tau) = \frac{1}{4} \sum_{(n,m) \in \mathbb{Z}^2} (n - im)^4 q^{n^2 + m^2}, \quad (8)$$

where as usual $q = e^{2\pi i\tau}$.

Theorem 2. *Formula (6) is true.*

Proof. Integrating (8) leads to

$$L(g, 4) = \frac{1}{4} \sum_{(n,m) \neq (0,0)} \frac{(n-im)^4}{(n^2+m^2)^4} = \frac{1}{4} \sum_{(n,m) \neq (0,0)} \frac{1}{(n+im)^4}.$$

The Weierstrass invariant $g_2(\tau)$ is defined by

$$g_2(\tau) := \frac{15}{4} \sum_{(n,m) \neq (0,0)} \frac{1}{(n+\tau m)^4},$$

and can be calculated using

$$g_2(\tau) = \frac{4}{3}(1-k^2+k^4)K^4(k), \quad (9)$$

where k, K and τ obey the classical relations $k = \theta_2^2(e^{\pi i \tau})/\theta_3^2(e^{\pi i \tau})$ and $K = \frac{\pi}{2}\theta_3^2(e^{\pi i \tau})$. (In the language of Eisenstein series, $g_2(\tau) = \frac{15}{2}\zeta(4)E_4(\tau)$.)

Standard evaluations show that $k = 1/\sqrt{2}$ when $\tau = i$, and thus

$$L(g, 4) = \frac{1}{15}g_2(i) = \frac{1}{15}K^4\left(\frac{1}{\sqrt{2}}\right).$$

Since $K(1/\sqrt{2}) = \Gamma^2(1/4)/(4\sqrt{\pi})$, we obtain

$$L(g, 4) = \frac{\Gamma^8(1/4)}{3840\pi^2},$$

completing the proof. □

Remark 1. Wan used numerical experiments to observe that the moments of K'^3 and K^2K' are related by a rational factor [14]:

$$\int_0^1 K'(k)^3 dk = 3 \int_0^1 K(k)^2 K'(k) dk. \quad (10)$$

We sketch a proof of equation (10) here. In particular, we can use Tricomi's Fourier series [14, Section 6]:

$$K(\sin t) = \sum_{n=0}^{\infty} \frac{\Gamma(n + \frac{1}{2})^2}{\Gamma(n+1)^2} \sin((4n+1)t),$$

to deduce that

$$\sum_{n=0}^{\infty} \frac{\Gamma^2(n + \frac{1}{2})}{\Gamma^2(n+1)} \int_0^{\pi/2} \cos((4n+2)t) \cos(t) F(\sin t) dt = \int_0^1 (k'K'(k) - kK(k)) F(k) dk, \quad (11)$$

whenever F is selected so that summation and integration are interchangeable. To derive (11), perform the change of variables $k \mapsto \cos t$ on the right hand side, and then use the trigonometric identity

$$\cos(t) \cos((4n+1)t) - \sin(t) \sin((4n+1)t) = \cos((4n+2)t).$$

We then set $F(k) = K^2(k)/k'$ in (11). The right-hand side simplifies under the transformation $k \mapsto k'$, and cancelations occur on the left-hand side due to orthogonality. After simplifying we have

$$\frac{\pi^2}{8} \sum_{n=0}^{\infty} \frac{\Gamma^4(n + \frac{1}{2})}{\Gamma^4(n + 1)} {}_4F_3 \left(\begin{matrix} \frac{1}{2}, \frac{1}{2}, -n, -n \\ 1, \frac{1}{2} - n, \frac{1}{2} - n \end{matrix}; 1 \right) = \int_0^1 K'(k)^3 - K(k)^2 K'(k) dk.$$

The sum is precisely $2 \int_0^1 K(k)^2 K'(k) dk$, as established near the end of [14]. This proves (10). \diamond

Before proving (5), we note that there are various reformulations of the integral. A quadratic transformation gives

$$\int_0^1 \frac{K'(k)^3}{\sqrt{k}(1-k^2)^{3/4}} dk = \frac{1}{8\sqrt{2}} \int_0^1 \frac{(1+k)^3 K'(k)^3}{k^{3/4} \sqrt{1-k}} dk.$$

Experimentally, we discover the binary theta expansion

$$h(\tau) = \prod_{n=1}^{\infty} \frac{(1 - (-1)^n q^{4n})^{14}}{(1 - q^{8n})^4} \quad (12)$$

$$= \frac{1}{2} \sum_{m,n=-\infty}^{\infty} (-1)^m (2n + 1 - 2im)^4 q^{(2m)^2 + (2n+1)^2}. \quad (13)$$

The proof of (13) is slightly tedious. We may show that both sides are modular forms on $\Gamma_1(64)$, that their Fourier coefficients agree for sufficiently many terms, and then appeal to the valence formula. Alternatively, the double sum can be built up from formulas of derivatives of $\theta_2(q^4)$ and $\theta_4(q^4)$ (Ramanujan's $\varphi(-q^4)$ and $\psi(q^8)$) in [2, pp. 122–123]. The result can be written in terms of k and K , and then reduced to the product (12).

Theorem 3. *Formula (5) is true.*

Proof. From (13), we have the (lattice) sum

$$2L(h, 4) = \sum_{m,n=-\infty}^{\infty} \frac{(-1)^m (2n + 1 - 2im)^4}{((2n + 1)^2 + (2m)^2)^4} = \sum_{m,n=-\infty}^{\infty} \frac{(-1)^m}{(2n + 1 + 2im)^4},$$

so we are reduced to showing

$$\sum_{m,n=-\infty}^{\infty} \frac{(-1)^m}{(2n + 1 + 2im)^4} = \frac{\Gamma^8(1/4)}{1024\sqrt{2}\pi^2}. \quad (14)$$

This can be achieved using routine manipulations on $g_2(\tau)$. By (9), we have

$$\sum_{(m,n) \neq (0,0)} \frac{(-1)^m}{(n + m\tau)^4} = \frac{4}{15} (2g_2(2\tau) - g_2(\tau)).$$

This leads to

$$\begin{aligned} \sum_{m,n} \frac{(-1)^m}{(2n+1+m\tau)^4} &= \frac{4}{15}(2g_2(2\tau) - g_2(\tau)) - \frac{1}{2^4} \frac{4}{15}(2g_2(\tau) - g_2(\tau/2)) \\ &= \frac{1}{60}(g_2(\tau/2) - 18g_2(\tau) + 32g_2(2\tau)). \end{aligned} \quad (15)$$

Now (14) follows by setting $\tau = 2i$ in (15) and simplifying. The required values of k (which are the 1st, 4th and 16th singular values) and K at these values can be found in [2] and [17]. \square

Example 1. We may look at other critical L -values of g and h . Simple calculations give

$$\begin{aligned} L(g, 3) &= \frac{1}{2\pi} \int_0^1 kK'(k)^2 K(k) dk, \\ L(g, 2) &= \frac{1}{\pi^2} \int_0^1 kK(k)^2 K'(k) dk, \\ L(g, 1) &= \frac{1}{\pi^3} \int_0^1 kK(k)^3 dk. \end{aligned}$$

One may notice the similarity between $L(g, 3)$ and $L(g, 2)$; indeed they are related by the substitution $k \mapsto k'$. This similarity can be explained, since g is a weight 5 modular form, so there is a *functional equation* connecting $L(g, s)$ and $L(g, 5-s)$ (with some Gamma factors). In view of (10), (2) and [14, Thm. 5], the values $\pi^{5-s}L(g, s)$ where $s \in \{1, 2, 3, 4\}$ are all related to each other by rational constants.

Similarly, we have $L(h, 1) = 384/\pi^3 L(h, 4)$, which can be explained by either using the functional equation or a substitution in the integral. Experimentally we also observe that $L(h, 4) = \frac{\pi}{4}L(h, 3)$; an integral formulation for this is

$$\int_0^1 \frac{K'(k)^2 K(k)}{\sqrt{k}(1-k^2)^{3/4}} dk = \frac{\Gamma^8(1/4)}{32\sqrt{2}\pi^2}, \quad (16)$$

and an equivalent formulation as a sum is

$$\sum_{m,n=-\infty}^{\infty} \frac{(-1)^m(2n+1-2mi)}{(2n+1+2im)^3} = \frac{\Gamma^8(1/4)}{256\sqrt{2}\pi^2}.$$

We are very grateful to Y. Zhou, who has kindly supplied us with a proof of (16) (plus some generalizations) using techniques from [15]; since such techniques differ significantly from the ones in the current paper, we do not include it here. Our approach to (16) is modular (using some ideas from [12]), but is substantially more complicated, and we will present it in a future paper.¹ \diamond

¹M. Rogers, J. G. Wan and I. J. Zucker, *Integrals from Lattice Sums*, work in progress, 2013.

Remark 2. Generalizations of (1) are possible. Start with the level 4, weight 9 newform

$$f(q) = \frac{1}{4} \sum_{m,n} (m - in)^8 q^{m^2+n^2},$$

then with $k = \theta_2^2(q)/\theta_3^2(q)$ as usual, it can be checked by computing the derivatives of θ_3 that

$$f(q) = \frac{8(4k^2k'^2 + k^4k'^4)}{\pi^9} K^9(k). \quad (17)$$

We have the following critical L -value,

$$L(f, 8) = \frac{1}{4} \sum_{(n,m) \neq (0,0)} \frac{1}{(m + in)^8} = \frac{\Gamma^{16}(\frac{1}{4})}{2^{10} 525 \pi^4}.$$

The last equality holds since the sum equals $\frac{1}{2}\zeta(8)E_8(i)$ and $E_8 = E_4^2$. Writing $L(f, 8)$ as an integral using (17), we obtain

$$\int_0^1 k(4 + k^2 - k^4)K'(k)^7 dk = \frac{3 \Gamma^{16}(\frac{1}{4})}{2^{12} 5 \pi^4}. \quad (18)$$

Experimentally, $L(f, 8 - i)/\pi^i$ are all related by rational constants. Since

$$\sum_{(n,m) \neq (0,0)} \frac{1}{(m + in)^{4k}}$$

can always be expressed as a rational number times a power of $K(1/\sqrt{2})$ [7], we have found generalizations of the above result which involve higher powers of K' , for instance

$$\int_0^1 k(16 - 92k^2 + 93k^4 - 2k^6 + k^8)K'(k)^{11} dk = \frac{189 \Gamma^{24}(\frac{1}{4})}{2^{15} 65 \pi^6}. \quad (19)$$

Moments of higher powers of K' , the observation about rational constants above, as well as many other integrals involving the complete elliptic integrals, will be elaborated in a future paper (see footnote 1). \diamond

4. WEIGHT 3 CASES AND LATTICE SUMS

In this section we note some additional formulas for critical L -values of weight 3 cusp forms. The ideas for the proof below are borrowed from [10].

Theorem 4. *Suppose that $f(\tau) = \eta^3(r\tau)\eta^3(s\tau)$, where $r + s \equiv 0 \pmod{8}$. Then*

$$L(f, 2) = \frac{8\pi^2}{\sqrt{rs^3}} \left(\sum_{n=0}^{\infty} (-1)^{\frac{n(n+1)}{2}} q^{\frac{(2n+1)^2}{8}} \right)^4 = \frac{8}{\sqrt{rs^3}} k_{r/s} k'_{r/s} K^2(k_{r/s}), \quad (20)$$

where $q = e^{-\pi\sqrt{r/s}}$, and k_p denotes the p th singular value of K .

Proof. The Jacobi triple product gives

$$\eta^3(\tau) = \sum_{n=1}^{\infty} n\chi_{-4}(n)e^{\frac{2\pi in^2\tau}{8}}.$$

If we write $L(f, 2)$ as a real-valued integral, then

$$L(f, 2) = 4\pi^2 \int_0^{\infty} \sum_{n,k \geq 1} nk\chi_{-4}(n)\chi_{-4}(k)e^{-\frac{2\pi rn^2u}{8} - \frac{2\pi sk^2u}{8}} u du.$$

Applying the involution for the weight $3/2$ theta function leads to

$$L(f, 2) = \frac{4\pi^2}{\sqrt{s^3}} \int_0^{\infty} \sum_{n,k \geq 1} nk\chi_{-4}(n)\chi_{-4}(k)e^{-\frac{2\pi rn^2u}{8} - \frac{2\pi k^2}{8su}} \frac{du}{\sqrt{u}}.$$

By absolute convergence, we evaluate the integral first using standard results – which produces a Bessel K function with order $1/2$. Simplifying, we get

$$\begin{aligned} L(f, 2) &= \frac{8\pi^2}{\sqrt{rs^3}} \sum_{n,k \geq 1} k\chi_{-4}(k)\chi_{-4}(n)q^{\frac{nk}{2}} = \frac{8\pi^2}{\sqrt{rs^3}} \sum_{k \geq 1} \frac{k\chi_{-4}(k)q^{\frac{k}{2}}}{1+q^k} \\ &= \frac{8\pi^2}{\sqrt{rs^3}} \left(\sum_{k=0}^{\infty} (-1)^{\frac{k(k+1)}{2}} q^{\frac{(2k+1)^2}{8}} \right)^4. \end{aligned}$$

The final Lambert series identity follows from results in Ramanujan's notebooks [2]. The connection with singular values follows from standard theta function manipulations, leading to

$$L(f, 2) = \frac{2\pi^2}{\sqrt{rs^3}} \theta_2^2(q)\theta_4^2(q),$$

and from the fact that $k_p = \theta_2^2(e^{-\pi\sqrt{p}})/\theta_3^2(e^{-\pi\sqrt{p}})$. □

Theorem 5. *The following evaluations are true:*

$f(\tau)$	$L(f, 2)$
$\eta^6(4\tau)$	$\frac{\Gamma^4(\frac{1}{4})}{64\pi}$
$\eta^3(2\tau)\eta^3(6\tau)$	$\frac{\Gamma^6(\frac{1}{3})}{2^{17/3}\pi^2}$
$\eta^3(\tau)\eta^3(7\tau)$	$\frac{\Gamma^2(\frac{1}{7})\Gamma^2(\frac{2}{7})\Gamma^2(\frac{4}{7})}{224\pi^2}$
$\eta^3(3\tau)\eta^3(5\tau) \pm \eta^3(\tau)\eta^3(15\tau)$	$\frac{\Gamma(\frac{1}{15})\Gamma(\frac{2}{15})\Gamma(\frac{4}{15})\Gamma(\frac{8}{15})}{30\sqrt{54} \mp 6\pi}$
$\frac{\eta^5(4\tau)\eta^5(8\tau)}{\eta^2(2\tau)\eta^2(16\tau)}$	$\frac{\Gamma^2(\frac{1}{8})\Gamma^2(\frac{3}{8})}{64\sqrt{2}\pi}$
$\frac{\eta^{18}(8\tau)}{\eta^6(4\tau)\eta^6(16\tau)}$	$\frac{\Gamma^4(\frac{1}{4})}{32\sqrt{2}\pi}$
$\eta^2(\tau)\eta(2\tau)\eta(4\tau)\eta^2(8\tau)$	$\frac{\Gamma^2(\frac{1}{8})\Gamma^2(\frac{3}{8})}{192\pi}$

Proof. For the first four entries in the table we use Theorem 4. The singular values k_p , as well as $K(k_p)$, are well tabulated in [2] and [17]. For instance, for the third entry we need

$$k_7 = \frac{\sqrt{2}(3 - \sqrt{7})}{8}, \quad K(k_7) = \frac{\Gamma(1/7)\Gamma(2/7)\Gamma(4/7)}{4\sqrt[4]{7}\pi}.$$

In general, k_p is algebraic and $K(k_p)$ involves only algebraic numbers and Gamma functions.

The last three entries are weight 3 cusp forms in Martin's list [9]. For the third last one, we convert it to the following integral,

$$L(f, 2) = \frac{1}{4\sqrt{2}} \int_0^1 \frac{K(k)}{\sqrt{k(1-k)}} dk,$$

which can be evaluated using the series representation of K and interchanging the order of integration and summation.

For the second last entry, we have

$$L(f, 2) = \frac{1}{8} \int_0^1 \frac{K(k)}{\sqrt{k(1-k^2)^{3/4}}} dk.$$

The evaluation of this integral follows from [14, equation (8)]. Indeed, many integrals over a single K have closed forms, and the two integrals we just evaluated can also be done by a computer algebra system.

For the last entry in the table, we obtain

$$L(f, 2) = \frac{1}{2\sqrt{2}} \int_0^1 \frac{K(k)}{\sqrt{1+k}} dk.$$

Denoting the integral by I_1 , we use the moments of K' found in [14] and a quadratic transformation to produce

$$I_1 = \frac{1}{\sqrt{2}} \int_0^1 \frac{K'(x)}{\sqrt{1+x}} dx = \frac{1}{2\sqrt{2}} \left[\frac{\Gamma^2(\frac{1}{8})\Gamma^2(\frac{3}{8})}{16\pi} - {}_4F_3\left(\frac{3}{4}, 1, 1, \frac{5}{4}; \frac{3}{2}, \frac{3}{2}, \frac{3}{2}; 1\right) \right]. \quad (21)$$

Similarly, with the auxiliary integral $I_2 := \int_0^1 \frac{K(x)}{\sqrt{x(1+x)}} dx$, we have

$$I_2 = \frac{1}{\sqrt{2}} \int_0^1 \frac{K'(x)}{\sqrt{1-x}} dx = \frac{1}{2\sqrt{2}} \left[\frac{\Gamma^2(\frac{1}{8})\Gamma^2(\frac{3}{8})}{16\pi} + {}_4F_3\left(\frac{3}{4}, 1, 1, \frac{5}{4}; \frac{3}{2}, \frac{3}{2}, \frac{3}{2}; 1\right) \right]. \quad (22)$$

Experimentally it is observed that $I_2 = 2I_1$, which can be shown as follows.

We know that $I_1/(2\sqrt{2}) = L(f, 2)$. On the other hand, it is readily verifiable that $I_2/(2\sqrt{2}) = L(f_0, 2)$, where $f_0(q) := -f(-q)$. Consequently, by looking at the q -expansion of $f_0 - f$ and using the fact that the coefficients of f are multiplicative, we deduce that $I_2 - I_1 = I_1$. The desired evaluation of I_1 now follows by combining (21) and (22), and inter alia we also obtain a closed form for the ${}_4F_3$. \square

Remark 3. Here we retain the notation used in the last part of the proof above. Firstly, with $f(\tau) = \eta^2(\tau)\eta(2\tau)\eta(4\tau)\eta^2(8\tau)$, it is true that

$$f(\tau) = \frac{1}{2} \sum_{(n,m) \neq (0,0)} (m^2 - 2n^2) q^{m^2+2n^2}.$$

Secondly, it is known that the series (3) for the L -value of a weight k cusp form converges conditionally for some $s < (k+1)/2$. In particular, the series for $L(f, 2)$ converges, and so by looking at the partial sums, we have

$$I_1 = \sqrt{2} \sum_{(m,n) \neq (0,0)} \frac{m^2 - 2n^2}{(m^2 + 2n^2)^2},$$

where we sum over expanding ellipses $m^2 + 2n^2 \leq M$, $M \rightarrow \infty$. Similarly,

$$I_2 = \sqrt{2} \sum_{(m,n) \neq (0,0)} (-1)^{m+1} \frac{m^2 - 2n^2}{(m^2 + 2n^2)^2}.$$

Subtracting the sums gives another proof that $I_2 - I_1 = I_1$. Since the sum for I_2 has better convergence properties (it can be summed over expanding rectangles), we will deal exclusively with alternating versions of the lattice sums we encounter. Note that we can decompose f into weight $3/2$ theta functions $(\eta^2(\tau)\eta^2(4\tau)/\eta(2\tau))$.

$(\eta^2(2\tau)\eta^2(8\tau)/\eta(4\tau))$ and therefore obtain a different series expansion. Such expansions imply the alternating sum identities

$$\sum_{(m,n) \neq (0,0)} \frac{(-1)^{m+1}(m^2 - 2n^2)}{(m^2 + 2n^2)^2} = \sum_{m,n} \frac{18(-1)^m(3m+1)(3n+1)}{((3m+1)^2 + 2(3n+1)^2)^2} = \frac{\Gamma^2(\frac{1}{8})\Gamma^2(\frac{3}{8})}{48\pi}. \quad (23)$$

◇

Example 2. The first entry in Theorem 5 ($f(\tau) = \eta^6(4\tau)$) corresponds to the result

$$L(f, 2) = \frac{1}{4} \int_0^1 \frac{K(k)}{\sqrt{1-k^2}} dk = \frac{\Gamma^4(\frac{1}{4})}{64\pi}, \quad (24)$$

and may be expressed as a lattice sum using a binary theta-expansion of f :

$$\sum_{(m,n) \neq (0,0)} (-1)^{m+1} \frac{m^2 - 4n^2}{(m^2 + 4n^2)^2} = \frac{\Gamma^4(\frac{1}{4})}{32\pi}. \quad (25)$$

Combined with well-known lattice sum evaluations [16], consequences of (25) include

$$\sum_{(m,n) \neq (0,0)} \frac{(-1)^m n^2}{(m^2 + 4n^2)^2} = \frac{\Gamma^4(\frac{1}{4})}{256\pi} - \frac{3\pi \log 2}{32}, \quad \sum_{(m,n) \neq (0,0)} \frac{(-1)^{m+1} m^2}{(m^2 + 4n^2)^2} = \frac{\Gamma^4(\frac{1}{4})}{64\pi} + \frac{3\pi \log 2}{8}.$$

On the other hand,

$$L(f, 1) = \frac{1}{\pi} \int_0^1 \frac{K(k)}{\sqrt{1-k^2}} dk = \frac{\Gamma^4(\frac{1}{4})}{16\pi^2},$$

which is consistent with the functional equation satisfied by $L(f, s)$.

We note that (24) actually is a specialization of [5, eqn after (5.14)], which states

$$\sum_{n=1}^{\infty} \frac{a_n}{n^2} q^{n/4} = \frac{\pi\sqrt{k}}{4K(k)} {}_3F_2\left(\begin{matrix} \frac{3}{4}, \frac{3}{4}, 1 \\ \frac{5}{4}, \frac{5}{4} \end{matrix}; k^2\right),$$

where a_n are the coefficients in the q -expansion of f , and as usual $k = \theta_2^2(q)/\theta_3^2(q)$; taking the limit as $k \rightarrow 1^-$ and appealing to the Stolz-Cesàro theorem recovers (24). ◇

Example 3. The second entry in Theorem 5 ($f(\tau) = \eta^3(2\tau)\eta^3(6\tau)$) gives the non-trivial integral evaluation

$$L(f, 2) = \int_0^1 (3+6p)^{-\frac{1}{2}} K\left(\frac{p^{\frac{3}{2}}(2+p)^{\frac{1}{2}}}{(1+2p)^{\frac{1}{2}}}\right) dp = \frac{\Gamma^6(\frac{1}{3})}{2^{\frac{17}{3}}\pi^2}, \quad (26)$$

where we have used the parametrization of the degree 3 modular equation and multiplier [2, Ch. 19]. We also have $L(f, 1) = \sqrt{3}/\pi L(f, 2)$, which can be shown from (26) either by the functional equation or by using a cubic transformation. We note that (26) (after a change of variable) appears in a very different context in [3, Sec. 3].

The lattice sum associated with (26) is

$$\sum_{(m,n) \neq (0,0)} (-1)^{m+n+1} \frac{m^2 - 3n^2}{(m^2 + 3n^2)^2} = \frac{\Gamma^6(\frac{1}{3})}{2^{\frac{14}{3}} \pi^2}. \quad (27)$$

For the third entry in Theorem 5, we do not seem to obtain a reasonably concise integral involving K (due to the apparent lack of a parametrization for the degree 7 modular equation). Using the binary theta function for $\eta^3(\tau)\eta^3(7\tau)$ [4], we obtain the sum

$$\sum_{(m,n) \neq (0,0)} \frac{(-1)^m (2n^2 - m^2)}{(m^2 + mn + 2n^2)^2} = \frac{\Gamma^2(\frac{1}{7})\Gamma^2(\frac{2}{7})\Gamma^2(\frac{4}{7})}{56\pi^2}. \quad (28)$$

◇

Example 4. An L -value of the function g used in the proof of (1) also gives some interesting lattice sum evaluations. Using the closed form for $L(g, 3)$, the multiplicativity of the coefficients of g , and results in [16], we deduce

$$\begin{aligned} \sum_{(m,n) \neq (0,0)} \frac{(-1)^{m+n} m^2 n^2}{(m^2 + n^2)^3} &= \frac{\Gamma^8(1/4)}{2^9 3 \pi^3} - \frac{\pi \log 2}{8}, \\ \sum_{(m,n) \neq (0,0)} \frac{(-1)^{m+n} m^4}{(m^2 + n^2)^3} &= -\frac{\Gamma^8(1/4)}{2^9 3 \pi^3} - \frac{3\pi \log 2}{8}, \\ \sum_{(m,n) \neq (0,0)} \frac{(-1)^m m^2 n^2}{(m^2 + n^2)^3} &= -\frac{\Gamma^8(1/4)}{2^{10} 3 \pi^3} - \frac{\pi \log 2}{16}. \end{aligned} \quad (29)$$

◇

5. EVEN WEIGHT CASES AND CONCLUSION

For *even weight* cusp forms we do not seem to obtain formulas for the L -values in terms of Gamma functions; instead hypergeometric functions are involved. Consider the following pair of weight 4 examples:

$$f_1(\tau) = \frac{\eta^{16}(4\tau)}{\eta^4(2\tau)\eta^4(8\tau)}, \quad f_2(\tau) = \eta^4(2\tau)\eta^4(4\tau).$$

All of the critical L -values of f_1 and f_2 reduce to special values of hypergeometric functions. Furthermore, there are some very curious relations *between* the L -values of both cusp forms:

$$\begin{aligned} L(f_1, 3) &= \frac{\pi}{2} L(f_2, 2) = \frac{\pi^2}{8} L(f_1, 1) = \frac{1}{8} \int_0^1 \frac{K(k)^2}{\sqrt{1-k^2}} dk \\ &= \frac{\pi^3}{32} {}_4F_3 \left(\begin{matrix} \frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \frac{1}{2} \\ 1, 1, 1 \end{matrix}; 1 \right). \end{aligned} \quad (30)$$

The last equality follows from [14, Eqn. 35]. Similarly we have

$$\begin{aligned} L(f_2, 3) &= \frac{\pi}{4}L(f_1, 2) = \frac{\pi^2}{4}L(f_2, 1) = \frac{1}{8} \int_0^1 \frac{K(k)K'(k)}{\sqrt{1-k^2}} dk = \frac{1}{4} \int_0^1 K(k)^2 dk \\ &= \frac{\pi^4}{128} {}_7F_6 \left(\begin{matrix} \frac{5}{4}, \frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \frac{1}{2} \\ \frac{1}{4}, 1, 1, 1, 1, 1 \end{matrix}; 1 \right). \end{aligned} \quad (31)$$

The second last equality comes from [14, Sec. 5.4], while the hypergeometric evaluation follows from [14, Eqns. 3 and 18]. Curiously this last integral also appears in connection with random walks [3, Sec. 3].

We conclude with the following observations and directions for further research.

- (1) It seems that the critical L -values of some even weight cusp forms can be expressed as hypergeometric functions, while those of odd weight cusp forms can often be expressed in terms of Gamma functions. It would be interesting to explain this discrepancy, because the Ramanujan zeta function, and various other interesting zeta functions, are attached to such cusp forms (see e.g. [11]). It would also be illuminating to see if many other critical L -values of odd weight cusp forms evaluate in terms of Gamma functions.
- (2) Some general results about the critical L -values of even weight cusp forms are recorded in [8]. For instance, as mentioned in section 3.4 there, the critical L -values satisfy an ‘alternating’ property, namely the ratios $L(f, 2k)/L(f, 2k-2)$ and $L(f, 2k+1)/L(f, 2k-1)$ can all be expressed in terms of algebraic numbers and Gamma functions (where s is the weight and $k = 1, 2, \dots$). It was simple to verify this for f_1 and f_2 studied above. On the other hand, in the two weight 5 cusp forms we studied, all the critical L -values are related by rational multiples and powers of π . (This also holds true numerically for the newform (17)). This phenomenon may have ultimately come from the functional equation for $L(f, s)$. For odd weight cusp forms (or newforms), how generally it is true that $L(f, s)/L(f, s-1)$ is an algebraic number times some Gamma factors?
- (3) As noted above, $L(f_1, s)/L(f_2, s-1)/\pi$ is a rational number. (Note that the q -expansions of f_1 and f_2 are also closely related.) Are there other pairs of even weight modular forms with the same property? If so, is there a method to find, given one function in the pair, the other function?

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